

## The Modified Tanh Method for Solving the Improved Eckhaus equation and the (2+1)-dimensional Improved Eckhaus Equation

N. Taghizadeh and M. Mirzazadeh

Department of Mathematics, Faculty of Science, University of Guilan, P.O.Box 1914, Rasht, Iran.

**Abstract:** The modified tanh method is one of most direct and effective algebraic method for obtaining exact solutions of nonlinear partial differential equations. The method can be applied to nonintegrable equations as well as to integrable ones. In this paper, we look for exact soliton solutions of the improved Eckhaus equation and the (2+1)-dimensional improved Eckhaus equation.  
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**Key words:** Modified tanh method; Improved Eckhaus equation; (2+ 1)-dimensional improved Eckhaus equation.

### INTRODUCTION

Nonlinear evolution equations have a major role in various scientific and engineering fields, such as, fluid mechanics, plasma physics, optical fibers, solid state physics, chemical kinematics, chemical physics and geochemistry. Nonlinear wave phenomena of dispersion, dissipation, diffusion, reaction and convection are very important in nonlinear wave equations. In recent years, quite a few methods for obtaining explicit traveling and solitary wave solutions of nonlinear evolution equations have been proposed. A variety of powerful methods, such as, tanh-sech method (Malfliet W., 1992; Khater A.H., 2002; Wazwaz AM, 2006), extended tanh method (El-Wakil S.A., 2007; Fan E., 2000; Wazwaz A.M., 2005), hyperbolic function method (Xia T.C., 2001), sine-cosine method (Wazwaz A.M., 2004; Yusufoglu E., 2006) Jacobi elliptic function expansion method (Inc M., 2005), F-expansion method (Sheng Zhang, 2006), and the first integral method (Feng Z.S., 2002; Ding TR, 1997). In recent years, there was interest in obtaining exact solutions of NLPDEs by the extended tanh method. The standard tanh method is developed by Malfliet (1992). Recently, Wazwaz investigated exact solutions of NLPDEs by the extended tanh method. Fan in presented the generalized tanh method for constructing the exact solutions of NLPDEs, such as, the (2 + 1)-dimensional sine-Gordon equation and the double sine-Gordon equation. The Eckhaus equation (M.J. Ablowitz, eads

$$i u_t + u_{xx} + 2(|u|^2)_{xx} u + |u|^4 u = 0, \quad u: \mathbb{R} \rightarrow \mathbb{C}.$$

This equation is of nonlinear Schrodinger type. Eckhaus was found in (Calogero, F., 1987) as an asymptotic multi-scale reduction of certain classes of nonlinear partial differential equations. In (Calogero, 1987), many of the properties of the Eckhaus equation were investigated. In (Calogero, 1987), the Eckhaus equation was (exactly) linearised by a change (dependent) of variable. The improved Eckhaus equation, and the (2+ 1)-dimensional improved Eckhaus equation are given by

$$i u_t + u_{xx} + 2(|u|^2)_{xx} u + |u|^4 u = 0,$$

and

$$i u_t + u_{xx} - u_{yy} + 2(|u|^2)_{xx} u + |u|^4 u = 0.$$

The aim of this paper is to find exact solutions of the improved Eckhaus equation, and the (2+ 1)-dimensional improved Eckhaus equation by the modified extended tanh method.

### 2. Modified Extended Tanh Method:

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**Corresponding Author:** N. Taghizadeh, Department of Mathematics, Faculty of Science, University of Guilan, P.O.Box 1914, Rasht, Iran.  
E-mail: taghizadeh@guilan.ac.ir

For given a nonlinear equation

$$F(u, u_x, u_y, u_t, u_{xx}, u_{xy}, u_{xt}, \dots) = 0, \tag{1}$$

when we look for its traveling wave solutions, the first step is to introduce the wave transformation u

$$(x,y,t) = U(\xi), \quad \xi = x + \gamma y + \lambda t \text{ and change Eq.(1) to an ordinary differential equation(ODE)}$$

$$H(U, U', U'', U''', \dots) = 0. \tag{2}$$

The next crucial step is to introduce a new variable  $\phi = \phi(\xi)$ , which is a solution of the Riccati equation

$$\frac{d\phi}{d\xi} = k + \phi^2 \tag{3}$$

The modified extended tanh method admits the use of the finite expansion:

$$u(x, y, t) = U(\xi) = \sum_{i=0}^m a_i \phi^i(\xi) + \sum_{i=1}^m b_i \phi^{-i}(\xi) \tag{4}$$

where the positive integer m is usually obtained by balancing the highest-order linear term with the nonlinear terms in Eq.(2). Expansion (4) reduces to the generalized tanh method [16] for  $b_i = 0, i = 1, \dots, m$ . Substituting Eq.(3) and Eq.(4) into Eq.(2) and then setting zero all coefficients of  $\phi^i(\xi)$ , we can obtain a system

of algebraic equations with respect to the constants  $k, \gamma, \lambda, a_0, \dots, a_m, b_1, \dots, b_m$ . Then we can determine the constants  $\gamma, \lambda, a_0, \dots, a_m, b_1, \dots, b_m$ . The Riccati equation (3) has the general solutions:

If  $k < 0$  then

$$\begin{aligned} \phi(\xi) &= -\sqrt{-k} \tanh(\sqrt{-k}\xi), \\ \phi(\xi) &= -\sqrt{-k} \coth(\sqrt{-k}\xi). \end{aligned} \tag{5}$$

If  $k = 0$  then

$$\phi(\xi) = -\frac{1}{\xi} \tag{6}$$

If  $k > 0$  then

$$\begin{aligned} \phi(\xi) &= -\sqrt{k} \tan(\sqrt{k}\xi), \\ \phi(\xi) &= -\sqrt{k} \cot(\sqrt{k}\xi). \end{aligned} \tag{7}$$

Therefore, by the sign test of k, we get exact soliton solutions of Eq. (1).

### 3 Improved Eckhaus Equation:

Let us consider the improved Eckhaus equation

$$iu_t + u_{xx} + 2(|u|^2)_{xx} u + |u|^4 u = 0 \tag{8}$$

where  $u = u(x,t)$ .

We introduce the wave transformation

$$u(x, t) = e^{i\theta} U(\varepsilon), \theta = \alpha x + \beta t, \xi = x + \lambda t \tag{9}$$

where  $\alpha, \beta$  and  $\lambda$  are real constants and  $U(\xi)$  is real function. Substituting (9) into Eq.(8) we obtain the relation

$$-(\beta + \alpha^2)U + U'' + 2(U')^2U + U^5 = 0$$

hence

$$-(\beta + \alpha^2)U + U'' + 4(U')^2U + 4U''U^2 + U^5 = 0. \tag{10}$$

By balancing the term  $U''$  with the term  $U^5$  in Eq.(10), we obtain  $m+2+2m=5m$ ,

$$m + 2 + 2m = 5m,$$

then  $m = 1$ .

The modified extended tanh method (4) admits the use of the finite expansion:

$$U(\xi) = a_0 + a_1\phi(\xi) + \frac{b_1}{\phi(\xi)}. \tag{11}$$

Substituting (11) into Eq.(10), making use of Eq.(3), and by using Maple, we get a system of algebraic equations, for  $a_0, a_1, b_1, \alpha, \beta$  and  $k$  in the form

$$\begin{aligned} \phi^5 : 12a_1^3 + a_1^5 &= 0, \\ \phi^4 : 20a_0a_1^2 + 5a_0a_1^4 &= 0, \\ \phi^3 : 5b_1a_1^4 + 16ka_1^3 + 10a_0^2a_1^3 + 12a_1^2b_1 + 8a_0^2a_1 + 2a_1 &= 0, \\ \phi^2 : 20b_1a_0a_1^3 + 24ka_0a_1^2 + 10a_0^3a_1^2 + 8b_1a_0a_1 &= 0, \\ \phi^1 : 4k^2a_1^3 + 10b_1^2a_1^3 + 16kb_1a_1^2 + 30b_1a_0^2a_1^2 + 4b_1^2a_1 \\ &\quad - (\beta + \alpha^2)a_1 + 5a_0^4a_1 + 8ka_0^2a_1 + 2ka_1 = 0, \\ \phi^0 : 30b_1^2a_0a_1^2 + 4k^2a_0a_1^2 + 20b_1a_0^3a_1 + 16ka_0b_1a_1 \\ &\quad - (\beta + \alpha^2)a_0 + a_0^5 + 4b_1^2a_0 = 0, \\ \phi^{-1} : 10b_1^3a_1^2 + 4k^2b_1a_1^2 + 16kb_1^2a_1 + 30b_1^2a_0^2a_1 \\ &\quad + 5b_1a_0^4 + 8kb_1a_0^2 + 2kb_1 + 4b_1^3 - (\beta + \alpha^2)b_1 = 0, \\ \phi^{-2} : 8k^2b_1a_0a_1 + 20b_1^3a_0a_1 + 24kb_1^2a_0 + 10b_1^2a_0^3 &= 0, \\ \phi^{-3} : 12k^2b_1^2a_1 + 5b_1^4a_1 + 8k^2b_1a_0^2 + 16kb_1^3 + 2k^2b_1 + 10b_1^3a_0^2 &= 0, \end{aligned}$$

**4 (2+1)-Dimensional Improved Eckhaus Equation:**

$$\begin{aligned} \phi^{-4} : 5b_1^4a_0 + 20k^2b_1^2a_0 &= 0, \\ \phi^{-5} : 12k^2b_1^3 + b_1^5 &= 0. \end{aligned}$$

These equations lead to the following cases:

$$k = \frac{1}{96}, \quad \beta = \frac{1}{64} - \alpha^2, \quad a_0 = a_1 = 0, \quad b_1 = \pm i \frac{\sqrt{3}}{48}, \tag{12}$$

$$k = \frac{1}{96}, \quad \beta = \frac{1}{64} - \alpha^2, \quad a_0 = 0, \quad a_1 = \pm 2i\sqrt{3}, \quad b_1 = 0, \tag{13}$$

$$k = \frac{1}{384}, \quad \beta = \frac{1}{64} - \alpha^2, \quad a_0 = 0, \quad a_1 = \pm 2i\sqrt{3}, \quad b_1 = \pm i \frac{\sqrt{3}}{192}, \tag{14}$$

$$k = -\frac{1}{192}, \quad \beta = \frac{1}{48} - \alpha^2, \quad a_0 = 0, \quad a_1 = \pm 2i\sqrt{3}, \quad b_1 = \pm i \frac{\sqrt{3}}{96}, \tag{15}$$

where  $\alpha$  is an arbitrary constant.

The sets (12) - (15) give the soliton solutions

$$u_1(x, t) = \pm \frac{i}{2\sqrt{2}} e^{i(\alpha x + (\frac{1-64\alpha^2}{64})t)} \cot\left(\frac{1}{4\sqrt{6}}(x - 2\alpha t)\right).$$

$$u_2(x, t) = \pm \frac{i}{2\sqrt{2}} e^{i(\alpha x + (\frac{1-64\alpha^2}{64})t)} \tan\left(\frac{1}{4\sqrt{6}}(x - 2\alpha t)\right).$$

$$u_3(x, t) = \pm \frac{i}{4\sqrt{2}} e^{i(\alpha x + (\frac{1-64\alpha^2}{64})t)} \left\{ \tan\left(\frac{1}{8\sqrt{6}}(x - 2\alpha t)\right) + \cot\left(\frac{1}{8\sqrt{6}}(x - 2\alpha t)\right) \right\}.$$

$$u_4(x, t) = \pm \frac{i}{4} e^{i(\alpha x + (\frac{1-48\alpha^2}{48})t)} \left\{ \tanh\left(\frac{1}{8\sqrt{3}}(x - 2\alpha t)\right) + \coth\left(\frac{1}{8\sqrt{3}}(x - 2\alpha t)\right) \right\}.$$

**4 (2+1)-Dimensional Improved Eckhaus Equation:**

In this section we study the (2+1)-dimensional improved Eckhaus equation

$$i u_t + u_{xx} - u_{yy} + 2(|u|^2)_{xx} u + |u|^4 u = 0. \tag{16}$$

where  $u = u(x, y, t)$ .

Let us consider the traveling wave solutions

$$u(x, y, t) = e^{i\theta} U(\xi), \quad \theta = \alpha x + \beta y + \delta t, \quad \xi = x + \gamma y + \lambda t, \tag{17}$$

where  $\alpha, \beta, \delta, \gamma, \lambda$  are real constants and  $U(\xi)$  is real function. Substituting (17) into Eq.(16), we find the relation  $\lambda = -2(\alpha - \beta \gamma)$ , then (16) is following nonlinear ordinary differential equation

$$(\beta^2 - \alpha^2 - \delta)U + (1 - \gamma^2)U'' + 2(U^2)''U + U^5 = 0. \tag{18}$$

**4 (2+1)-Dimensional Improved Eckhaus Equation:**

It is easy to show that  $m = 1$ , if balancing  $U^5$  with  $U''$ .

The modified extended tanh method (4) admits the use of the finite expansion:

$$U(\xi) = a_0 + a_1\phi(\xi) + \frac{b_1}{\phi(\xi)}. \tag{19}$$

Substituting Eqs.(3) and (19) into Eq.(18), and equating the coefficients of like powers of  $\phi$  ( $i = -5, -4, -3, -2, -1, 0, 1, 2, 3, 4, 5$ ) to zero yields the system of algebraic equations to  $a_0, a_1, b_1, \alpha, \delta, \beta, \gamma$  and  $k$

$$\begin{aligned} \phi^5 : 12a_1^3 + a_1^5 &= 0, \\ \phi^4 : 20a_0a_1^2 + 5a_0a_1^4 &= 0, \\ \phi^3 : 5b_1a_1^4 + 16ka_1^3 + 10a_0^2a_1^3 + 12a_1^2b_1 + 8a_0^2a_1 + 2(1 - \gamma^2)a_1 &= 0, \\ \phi^2 : 20b_1a_0a_1^3 + 24ka_0a_1^2 + 10a_0^3a_1^2 + 8b_1a_0a_1 &= 0, \\ \phi^1 : 4k^2a_1^3 + 10b_1^2a_1^3 + 16kb_1a_1^2 + 30b_1a_0^2a_1^2 + 4b_1^2a_1 \\ + (\beta^2 - \alpha^2 - \delta)a_1 + 5a_0^4a_1 + 8ka_0^2a_1 + 2(1 - \gamma^2)ka_1 &= 0, \\ \phi^0 : 30b_1^2a_0a_1^2 + 4k^2a_0a_1^2 + 20b_1a_0^3a_1 + 16ka_0b_1a_1 \\ + (\beta^2 - \alpha^2 - \delta)a_0 + a_0^5 + 4b_1^2a_0 &= 0, \\ \phi^{-1} : 10b_1^3a_1^2 + 4k^2b_1a_1^2 + 16kb_1^2a_1 + 30b_1^2a_0^2a_1 \\ + 5b_1a_0^4 + 8kb_1a_0^2 + 2(1 - \gamma^2)kb_1 + 4b_1^3 + (\beta^2 - \alpha^2 - \delta)b_1 &= 0, \\ \phi^{-2} : 8k^2b_1a_0a_1 + 20b_1^3a_0a_1 + 24kb_1^2a_0 + 10b_1^2a_0^3 &= 0, \\ \phi^{-3} : 12k^2b_1^2a_1 + 5b_1^4a_1 + 8k^2b_1a_0^2 + 16kb_1^3 + 2(1 - \gamma^2)k^2b_1 + 10b_1^3a_0^2 &= 0, \\ \phi^{-4} : 5b_1^4a_0 + 20k^2b_1^2a_0 &= 0, \\ \phi^{-5} : 12k^2b_1^3 + b_1^5 &= 0. \end{aligned}$$

These equations lead to the following cases:

**Case A:**

$$k = \frac{1}{96}(1 - \gamma^2), \quad \delta = \beta^2 - \alpha^2 + \frac{1}{64} - \frac{1}{32}\gamma^2 + \frac{1}{64}\gamma^4, \quad a_0 = 0, \quad a_1 = \pm 2i\sqrt{3}, \quad b_1 = 0,$$

with  $\alpha, \beta, \gamma$  being arbitrary constants. By using (6) and (7) and (5) we obtain three solutions, namely,

$$u_1(x, y, t) = \pm \frac{2\sqrt{3}i}{x \pm y - 2(\alpha \pm \beta)t} e^{i(\alpha x + \beta y + (\beta^2 - \alpha^2)t)},$$

for  $k = 0$ .

$$u_2(x, y, t) = \pm \frac{i\sqrt{1 - \gamma^2}}{2\sqrt{2}} e^{i(\alpha x + \beta y + (\beta^2 - \alpha^2 + \frac{1}{64} - \frac{1}{32}\gamma^2 + \frac{1}{64}\gamma^4)t)} \tan\left(\frac{\sqrt{1 - \gamma^2}}{4\sqrt{6}}(x + \gamma y - 2(\alpha - \beta\gamma)t)\right),$$

for  $k > 0$ .

$$u_3(x, y, t) = \pm \frac{i\sqrt{\gamma^2 - 1}}{2\sqrt{2}} e^{i(\alpha x + \beta y + (\beta^2 - \alpha^2 + \frac{1}{64} - \frac{1}{32}\gamma^2 + \frac{1}{64}\gamma^4)t)} \tanh\left(\frac{\sqrt{\gamma^2 - 1}}{4\sqrt{6}}(x + \gamma y - 2(\alpha - \beta\gamma)t)\right),$$

for  $k < 0$ . Case B:

$$k = \frac{1}{96}(1 - \gamma^2), \quad \delta = \beta^2 - \alpha^2 + \frac{1}{64} - \frac{1}{32}\gamma^2 + \frac{1}{64}\gamma^4, \quad a_0 = a_1 = 0, \quad b_1 = \pm i \frac{\sqrt{3}}{48}(1 - \gamma^2),$$

with  $\alpha, \beta, \gamma$  being arbitrary constants. By using (7) and (5) we obtain two solutions, namely,

$$u_4(x, y, t) = \pm \frac{i\sqrt{1-\gamma^2}}{2\sqrt{2}} e^{i(\alpha x + \beta y + (\beta^2 - \alpha^2 + \frac{1}{64} - \frac{1}{32}\gamma^2 + \frac{1}{64}\gamma^4)t)} \cot\left(\frac{\sqrt{1-\gamma^2}}{4\sqrt{6}}(x + \gamma y - 2(\alpha - \beta\gamma)t)\right),$$

for  $k > 0$ .

$$u_5(x, y, t) = \pm \frac{i\sqrt{\gamma^2-1}}{2\sqrt{2}} e^{i(\alpha x + \beta y + (\beta^2 - \alpha^2 + \frac{1}{64} - \frac{1}{32}\gamma^2 + \frac{1}{64}\gamma^4)t)} \coth\left(\frac{\sqrt{\gamma^2-1}}{4\sqrt{6}}(x + \gamma y - 2(\alpha - \beta\gamma)t)\right),$$

for  $k < 0$ .

### 5. Conclusion:

In this paper, the modified tanh method has been successfully applied to find the solution for two nonlinear partial differential equations such as improved Eckhaus and (2 + 1)-dimensional improved Eckhaus equations. Thus, we can say that the proposed method can be extended to solve the problems of nonlinear partial differential equations which arising in the theory of solitons and other areas.

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